

For a state for  $l=0$ ,

$$\psi(\alpha) \rightarrow \psi^*(\alpha)$$

$$\mathbb{1}|\alpha\rangle = \mathbb{1} \int dx |x\rangle \langle x|\alpha\rangle = \int dx |x\rangle \langle x|\alpha\rangle^*$$

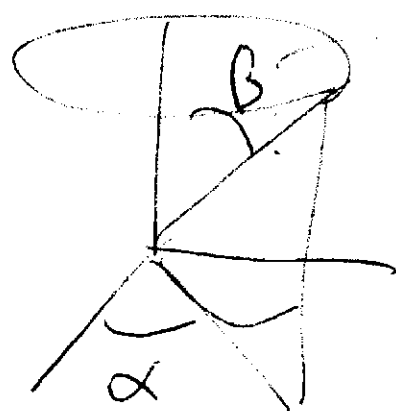
if we expand in terms of the momentum eigenstate,

$$\mathbb{1}|\alpha\rangle = \mathbb{1} \int dp |p\rangle \langle p|\alpha\rangle = \int dp |-p\rangle \langle p|\alpha\rangle^* \\ \phi(p) \rightarrow \underline{\underline{\phi^*(-p)}}$$

Time reversal for a spin  $\frac{1}{2}$  system

Sec. 3.2

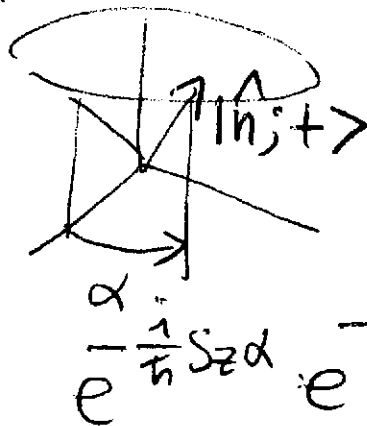
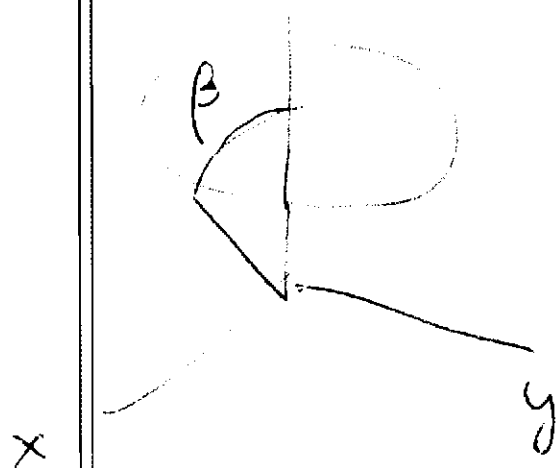
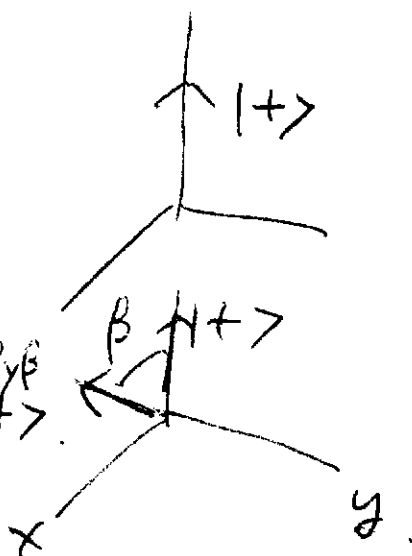
$|\hat{n}\rangle$



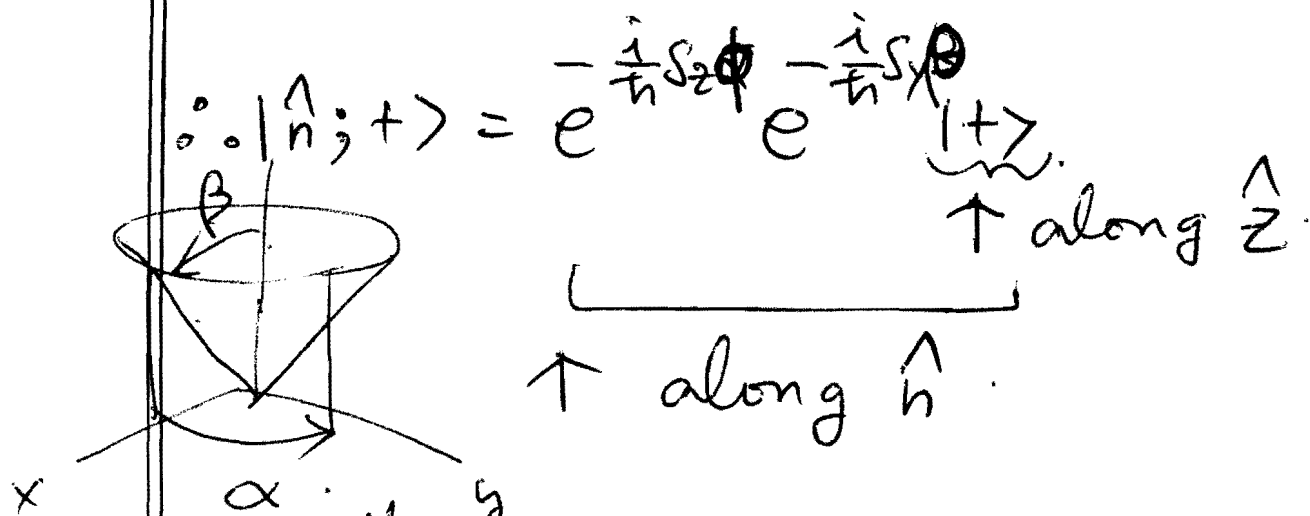
polar angle.

azimuthal angle

$$e^{-\frac{i}{\hbar} S_y \beta} |+\rangle$$



$$e^{-\frac{i}{\hbar} S_z \alpha} e^{-\frac{i}{\hbar} S_y \beta} |+\rangle$$



applying the time-reversal operator  $\Theta$ ,

$$\Theta |\hat{n}; +\rangle = e^{-\frac{i}{\hbar} S_z \theta} e^{-\frac{i}{\hbar} S_y \theta} \Theta |+\rangle$$

note that  $\Theta$

$$\Theta S_z \Theta^{-1} = -S_z \quad \rightarrow \quad \Theta e^{-\frac{i}{\hbar} S_z \theta} \Theta^{-1} = e^{\frac{i}{\hbar} S_z \theta}$$

$$\Theta S_y \Theta^{-1} = -S_y \quad \rightarrow \quad \Theta e^{-\frac{i}{\hbar} S_y \theta} \Theta^{-1} = e^{\frac{i}{\hbar} S_y \theta}$$

$$\Theta i \Theta^{-1} = -i$$

$$\Theta |\hat{n}; +\rangle = \eta |\hat{n}; -\rangle$$

$$|\hat{n}; -\rangle = e^{-\frac{i}{\hbar} S_z \theta} e^{-\frac{i}{\hbar} S_y (\pi + \theta)} |+\rangle$$

$$\eta |\hat{n}; -\rangle = \eta e^{-\frac{i}{\hbar} S_z \theta} e^{-\frac{i}{\hbar} S_y \theta} e^{-\frac{i}{\hbar} S_y \pi} |+\rangle$$

$$= e^{-\frac{i}{\hbar} S_z \theta} e^{-\frac{i}{\hbar} S_y \theta} \Theta |+\rangle$$

rotat  $e^{-\frac{i}{\hbar} S_y \pi} |+\rangle$   
 "  $K |+\rangle$   
 Because  $K |+\rangle = |+\rangle$

$$\Theta \text{ (circled) } = \eta e^{-\frac{i}{\hbar} S_y \pi} \text{ (circled) } K \leftarrow \text{complex conjugate op.}$$

unitary part U

$$\textcircled{H} = \eta e^{-\frac{i}{\hbar} S_y \pi}$$

$$e^{-\frac{i}{\hbar} S_y \pi} = e^{-\frac{i}{\hbar} \frac{\hbar}{2} \sigma_2 \pi} = e^{-\frac{i}{2} \pi \sigma_2}$$

$$e^{i \vec{\sigma} \cdot \vec{a}} = \sum_{k=0}^{\infty} i^{2k} \frac{(\vec{\sigma} \cdot \vec{a})^{2k}}{(2k)!} + \sum_{k=0}^{\infty} i^{2k+1} \frac{(\vec{\sigma} \cdot \vec{a})^{2k+1}}{(2k+1)!}$$

$$= \sum_{k=0}^{\infty} (-1)^k \frac{a^{2k}}{(2k)!} + i \vec{\sigma} \cdot \vec{a} \sum_{k=0}^{\infty} \frac{a^{2k+1}}{(2k+1)!} (-1)^k$$

$$(\vec{\sigma} \cdot \vec{a})^2 = \sigma_i \sigma_j a_i a_j = \frac{1}{2} (\sigma_i \sigma_j + \sigma_j \sigma_i) a_i a_j$$

$$= \delta_{ij} a_i a_j = a^2$$

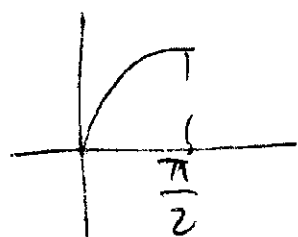
$$= \mathbb{1} \cos a + i \vec{\sigma} \cdot \vec{a} \sin a$$

$$e^{-\frac{i}{\hbar} S_y \pi} = e^{-\frac{i}{2} \pi \sigma_2}$$

$$= \mathbb{1} \cos\left(-\frac{\pi}{2}\right) + i \sigma_2 \sin\left(-\frac{\pi}{2}\right)$$

$$= -i \sigma_2$$

$$\sigma_2 = \frac{2}{\hbar} S_y$$



$$\sin\left(-\frac{\pi}{2}\right) = -1$$

$$\cos\left(-\frac{\pi}{2}\right) = 0$$

$$\therefore \textcircled{H} = \eta e^{-\frac{i}{\hbar} S_y \pi} = -\eta i \sigma_2 = -i \eta \left(\frac{2S_y}{\hbar}\right)$$

↑  
arbitrary phase

$$S_y = \frac{\hbar}{2} \sigma_2$$

$$\textcircled{H} = -i\eta \left( \frac{2S_y}{\hbar} \right) K = -i\eta \sigma_2 K$$

$$|+\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad S_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$e^{i\vec{\sigma} \cdot \vec{a}} = \mathbb{1} \cos a + i \vec{\sigma} \cdot \hat{a} \sin a$$

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\sigma_3 \sigma_1 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = i \sigma_2$$

$$\begin{aligned} |\hat{n}; +\rangle &= e^{-\frac{i}{\hbar} S_z \phi} e^{-\frac{i}{\hbar} S_y \theta} \\ &= e^{-\frac{i}{\hbar} \frac{\hbar}{2} \sigma_3 \phi} e^{-\frac{i}{\hbar} \frac{\hbar}{2} \sigma_2 \theta} \\ &= e^{+i\sigma_3(-\frac{\phi}{2})} e^{+i\sigma_2(-\frac{\theta}{2})} \end{aligned}$$

$$\begin{aligned} e^{i\sigma_3(-\frac{\phi}{2})} &= \mathbb{1} \cos(-\frac{\phi}{2}) + i\sigma_3 \sin(-\frac{\phi}{2}) \\ &= \mathbb{1} \cos \frac{\phi}{2} + \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} (-\sin \frac{\phi}{2}) \\ &= \begin{pmatrix} \cos \frac{\phi}{2} - i \sin \frac{\phi}{2} & 0 \\ 0 & \cos \frac{\phi}{2} + i \sin \frac{\phi}{2} \end{pmatrix} = \begin{pmatrix} e^{-\frac{i}{2}\phi} & 0 \\ 0 & e^{\frac{i}{2}\phi} \end{pmatrix} \end{aligned}$$

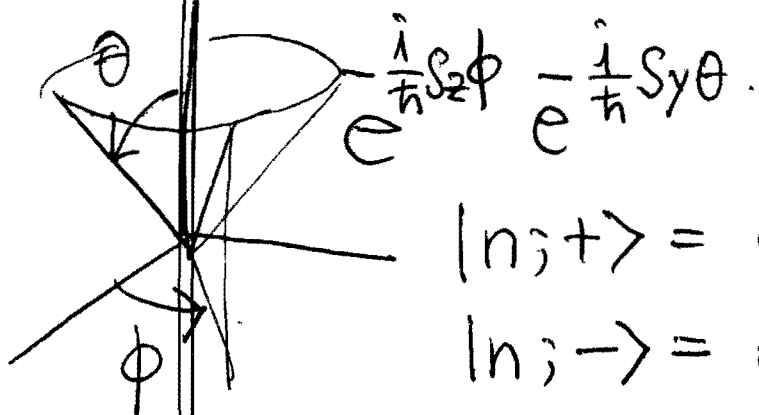
$$\begin{aligned} e^{i\sigma_2(-\frac{\theta}{2})} &= \mathbb{1} \cos(-\frac{\theta}{2}) + i\sigma_2 \sin(-\frac{\theta}{2}) \\ &= \begin{pmatrix} \cos \frac{\theta}{2} & 0 \\ 0 & \cos \frac{\theta}{2} \end{pmatrix} + \begin{pmatrix} 0 & -\sin \frac{\theta}{2} \\ \sin \frac{\theta}{2} & 0 \end{pmatrix} \\ &= \begin{pmatrix} \cos \frac{\theta}{2} & -\sin \frac{\theta}{2} \\ \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix} \end{aligned}$$

$$-i\sigma_2 = -i \begin{pmatrix} i & -i \\ i & -i \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$$

$$\begin{aligned} & \therefore e^{i\sigma_3(-\frac{\phi}{2})} e^{i\sigma_2(-\frac{\theta}{2})} \\ &= \begin{pmatrix} e^{-i\frac{\phi}{2}} & 0 \\ 0 & e^{i\frac{\phi}{2}} \end{pmatrix} \begin{pmatrix} \cos\frac{\theta}{2} & -\sin\frac{\theta}{2} \\ \sin\frac{\theta}{2} & \cos\frac{\theta}{2} \end{pmatrix} \\ &= \begin{pmatrix} e^{-i\frac{\phi}{2}} \cos\frac{\theta}{2} & -e^{-i\frac{\phi}{2}} \sin\frac{\theta}{2} \\ e^{i\frac{\phi}{2}} \sin\frac{\theta}{2} & e^{i\frac{\phi}{2}} \cos\frac{\theta}{2} \end{pmatrix} \end{aligned}$$

$$\begin{aligned} \therefore |\hat{n}; +\rangle &= \begin{pmatrix} \phantom{0} \\ \phantom{0} \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} e^{-i\frac{\phi}{2}} \cos\frac{\theta}{2} \\ e^{i\frac{\phi}{2}} \sin\frac{\theta}{2} \end{pmatrix} \\ |\hat{n}; -\rangle &= \begin{pmatrix} \phantom{0} \\ \phantom{0} \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} -e^{-i\frac{\phi}{2}} \sin\frac{\theta}{2} \\ e^{i\frac{\phi}{2}} \cos\frac{\theta}{2} \end{pmatrix} \end{aligned}$$

$$\begin{aligned} \textcircled{H} |\hat{n}; \ominus\rangle &= -i\eta \sigma_2 K |\hat{n}; +\rangle \quad -i\sigma_2 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \text{ same } \rho_0 \\ &= \eta \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \textcircled{K} \begin{pmatrix} e^{-i\frac{\phi}{2}} \cos\frac{\theta}{2} \\ e^{i\frac{\phi}{2}} \sin\frac{\theta}{2} \end{pmatrix} \quad \text{Complex conjugate} \\ &= \eta \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} e^{i\frac{\phi}{2}} \cos\frac{\theta}{2} \\ -e^{-i\frac{\phi}{2}} \sin\frac{\theta}{2} \end{pmatrix} = \eta \begin{pmatrix} -e^{-i\frac{\phi}{2}} \sin\frac{\theta}{2} \\ e^{i\frac{\phi}{2}} \cos\frac{\theta}{2} \end{pmatrix} \end{aligned}$$



$$e^{\frac{i}{\hbar} S_z \phi} e^{-\frac{i}{\hbar} S_y \theta}$$

$$|n; +\rangle = e^{-\frac{i}{\hbar} S_z \phi} e^{-\frac{i}{\hbar} S_y \theta} |+\rangle$$

$$|n; -\rangle = e^{-\frac{i}{\hbar} S_z \phi} e^{-\frac{i}{\hbar} S_y (\theta + \pi)} |+\rangle$$

$$\begin{pmatrix} 1 \\ 0 \end{pmatrix} \longrightarrow \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$|+\rangle \longrightarrow |-\rangle$$

$$e^{-\frac{i}{\hbar} S_y \pi} = e^{-\frac{i}{\hbar} \frac{\hbar}{2} \sigma_2 \pi} = e^{i \sigma_2 (-\frac{\pi}{2})} \quad i \sigma_2 = i \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$$

$$= \mathbb{1} \cos(-\frac{\pi}{2}) + i \sigma_2 \sin(-\frac{\pi}{2})$$

$$= \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \cos \frac{\pi}{2} + \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \sin \frac{\pi}{2}$$

$$\cos \frac{\pi}{2} = 0, \sin \frac{\pi}{2} = 1$$

$$= \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$$

$$\therefore \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$\textcircled{H} |n; +\rangle = \eta |n; -\rangle$$

$$\textcircled{H} e^{-\frac{i}{\hbar} S_z \phi} e^{-\frac{i}{\hbar} S_y \theta} |+\rangle = \eta e^{-\frac{i}{\hbar} S_z \phi} e^{-\frac{i}{\hbar} S_y (\theta + \pi)} |+\rangle$$

$$\textcircled{H} e^{-\frac{i}{\hbar} S_z \phi} \textcircled{H}^\dagger = e^{-\frac{(-i)}{\hbar} (-S_z) \phi} = e^{-\frac{i}{\hbar} S_z \phi}$$

$$\dots \textcircled{H} |+\rangle = \eta e^{-\frac{i}{\hbar} S_y \pi} |+\rangle = \eta e^{-\frac{i}{\hbar} S_y \pi} K |+\rangle$$

Because  $K |+\rangle = |+\rangle$  ( $S_z = \frac{\hbar}{2} (1 -)$ ) real  $|+\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ ,  $|-\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ .



$$\sqrt{(c^+|+\rangle + c^-|-\rangle)}$$

$$\textcircled{H}^2 = \eta \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} K \eta \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} K \begin{pmatrix} c^+ \\ c^- \end{pmatrix}$$

$$(z^*)^* = z$$

$$\textcircled{H}^2 = \eta \eta^* \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} K^2 \begin{pmatrix} c^+ \\ c^- \end{pmatrix}$$

$$= |\eta|^2 \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} c^+ \\ c^- \end{pmatrix}$$

$$|\eta|^2 = 1$$

unit modulus.

$$= - \begin{pmatrix} c^+ \\ c^- \end{pmatrix}$$

$$= - (c^+|+\rangle + c^-|-\rangle)$$

$$\therefore \textcircled{H}^2 = -\underline{1}$$

independent of the phase  $\eta$ .

spin  $-\frac{1}{2}$  state

$\textcircled{H}^2 = +\underline{1}$  for spin 0 state

We now prove that

$$\textcircled{H}^2 |j = \text{half-integer}\rangle = - |j\rangle$$

$$\textcircled{H}^2 |j = \text{integer}\rangle = + |j\rangle$$

therefore, eigenvalue of  $\textcircled{H}^2$  is  $(-1)^{2j}$ .



We recall that, for a spin- $\frac{1}{2}$  system,

$$\hat{H} = \eta e^{-\frac{i}{\hbar} S_y \pi} K$$

In general,

$$\begin{cases} |\hat{n}; jm\rangle = e^{-\frac{i}{\hbar} J_z \phi} e^{-\frac{i}{\hbar} J_y \theta} |\hat{z}; jm\rangle \\ \hat{H} |\hat{n}; jm\rangle = \hat{H} e^{-\frac{i}{\hbar} J_z \phi} e^{-\frac{i}{\hbar} J_y \theta} |\hat{z}; jm\rangle \end{cases}$$

$$\begin{aligned} & \hat{H} i \hat{H}^\dagger = -i \\ & \hat{H} \vec{J} \hat{H}^\dagger = -\vec{J} \\ & \Rightarrow \hat{H} e^{-\frac{i}{\hbar} \vec{J} \cdot \vec{\phi}} \hat{H}^\dagger = e^{-\frac{i}{\hbar} \vec{J} \cdot \vec{\phi}} \\ & \downarrow \\ & = \hat{H} e^{-\frac{i}{\hbar} J_z \phi} e^{-\frac{i}{\hbar} J_y \theta} \hat{H} |\hat{z}; jm\rangle \\ & = \eta |\hat{z}; j, -m\rangle \\ & = \eta \left[ e^{-\frac{i}{\hbar} J_z \phi} e^{-\frac{i}{\hbar} J_y \theta} |\hat{z}; j, -m\rangle \right] \\ & = \eta |\hat{n}; j, -m\rangle \end{aligned}$$

On the other hand

$$|\hat{n}; j, -m\rangle = e^{-\frac{i}{\hbar} J_z \phi} e^{-\frac{i}{\hbar} J_y \theta} e^{\frac{i}{\hbar} J_y \pi} |\hat{z}; j, +m\rangle$$

$$= \eta e^{-\frac{i}{\hbar} J_z \phi} e^{-\frac{i}{\hbar} J_y \theta} \eta e^{-\frac{i}{\hbar} J_y \pi} |\hat{z}; j, +m\rangle$$

$\therefore \hat{H} = \eta e^{-\frac{i}{\hbar} J_y \pi} K$  for a system with complex conjugation operator  $J$ .

$$\textcircled{H} = \eta e^{-\frac{i}{\hbar} J_y \pi} K$$

$$|\alpha\rangle = \sum |jm\rangle \langle jm|\alpha\rangle$$

Note that  $J_y |jm\rangle \neq m\hbar |jm\rangle$   
 { not  $J_y |jm\rangle = m\hbar |jm\rangle$   
 { but  $J_z |jm\rangle = m\hbar |jm\rangle$

~~Applying~~  $\textcircled{H}^2$

$$\textcircled{H}^2 = \eta e^{-\frac{i}{\hbar} J_y \pi} K \eta e^{-\frac{i}{\hbar} J_y \pi} K$$

$$= \eta e^{-\frac{i}{\hbar} J_y \pi} \eta^* K e^{-\frac{i}{\hbar} J_y \pi} K$$

$$= |\eta|^2 e^{-\frac{i}{\hbar} J_y \pi} e^{+\frac{i}{\hbar} J_y^* \pi} K^2$$

$$K^2 = \mathbb{1}$$

$$= |\eta|^2 e^{-\frac{i}{\hbar} J_y \pi} e^{+\frac{i}{\hbar} J_y^* \pi} = |\eta|^2 e^{-\frac{i}{\hbar} J_y (2\pi)}$$

$$J_{\pm} = J_x \pm i J_y$$

$$J_y^* = -J_y$$

$$\therefore J_y = \frac{1}{2i} (J_+ - J_-) \text{ and } J_x = \frac{1}{2} (J_+ + J_-)$$

$$\langle j'm'| J_{\pm} |jm\rangle = \hbar \sqrt{j(j+1) - m(m\pm 1)} \delta_{m', m\pm 1}$$

Therefore,  $J_{\pm}$  are ~~para~~ real.

$\Rightarrow J_y$  is pure imaginary  
 and  $J_x$  is real for any  $J$ .

$$\text{Therefore, } J_y^* = -J_y$$

Therefore,  $(\hat{H})^2 = e^{-\frac{i}{\hbar} \hat{J}_y (2\pi)}$

This is the rotation about the  $y$  axis by an angle of  $2\pi$ .

Homework)

Show that  $e^{-\frac{i}{\hbar} \hat{J}_y (2\pi)} |j m\rangle = (-1)^{2j} |j m\rangle$

---

One example is  $\hat{J}_y = \frac{\hbar}{2} \sigma_2$

$$\begin{aligned} e^{-\frac{i}{\hbar} \hat{J}_y (2\pi)} &= e^{-\frac{i}{\hbar} \frac{\hbar}{2} \sigma_2 (2\pi)} = e^{i\sigma_2 (-\pi)} \\ &= \mathbb{1} \cos(-\pi) + i\sigma_2 \sin(-\pi) \\ &= \cos \pi \mathbb{1} = -\mathbb{1}. \end{aligned}$$